Darle Wile Suppose

Theresiden.

$$Y(p_{\theta}) = X(p_{\theta})X(p_{\theta}) = \{q + p_{\theta}(q) = p_{\theta}\}$$

Finding the partial fraction expansion of $Y(j\omega)$ and taking the inverse Fourier trans-

$$y = 0 = \frac{1}{2}e^{-2x}w = 0 = \frac{1}{2}e^{-x}w = 0$$

a.3th. (ad) From the given information,

$$|U(1)\omega(t)| = \frac{\sqrt{\omega^2 + \omega^2}}{\sqrt{\omega^2 + \omega^2}} = 1.$$

Citi Seek

$$-ich((p_m)) = -tan^{-1}\frac{m}{2} - tan^{-1}\frac{m}{2} = -2tan^{-1}\frac{m}{2}$$

St. Stead

$$M(|f_{w}|) = -1 + \frac{5u}{u + 1}$$
 or $A(0) = -f(0) + 2uv^{-ut}u(0)$

(26) If a - 2, we have

$$||\mu(r(p,\omega))|| = 1, \qquad \text{with } |(p,\omega)| = -2\pi n n^{-1} \omega.$$

Thursdoon,

$$\chi(t) = \min(\frac{1}{\sqrt{3}} - \frac{\pi}{3}) - \min(t - \frac{\pi}{3}) + \min(\sqrt{3}t - \frac{2\pi}{3})$$

s.26. (a) The Soquency response t

$$M(p_{i+1}) = \frac{V(p_{i+1})}{V(p_{i+1})} = \frac{2(2 + p_{i+1})}{(4 + p_{i+1})(2 + p_{i+1})}$$

(b) Finding the partial fraction expansion of narrow in part (a) and taking its soverage Former transferm, we obtain

OKS We have

$$\frac{V'(y,\omega)}{V'(y,\omega)} = \frac{(y+y)(\omega)}{y+y} + \frac{(y+y)(y-y)}{y+y}$$

Cross multiplying and taking the inverse Feorier transform, we obtain

180

a.34. (a) Applying a Sequency shift to the analysis equation, we have

$$X(|y|\omega - \omega_0)) = \int_{-\infty}^{\infty} y(y)e^{-jkx-\omega_0}dy = \int_{-\infty}^{\infty} y(y)e^{jk\omega_0}e^{-j\omega_0}dt = FT(y)(y)e^{jk\omega_0}$$

(Did) Why Bullets

Alberta

$$a(t)a(t) \stackrel{dT_0}{=} \frac{1}{3^n} (X(p_0)) \cdot W(p_0))$$

 $= X(p_0) \cdot \delta(p_0 - p_0)$
 $= X(p_0 - p_0)$

s.30. (a) From the Fourier transferor analysis equation, we have

$$G(p\omega) = \int_{-\infty}^{\infty} p(t)e^{-(pst)}dt = \int_{-\infty}^{\infty} X(pt)e^{-(pst)}dt.$$
 (5 = 10 - 1

Also from the France transform synthesis equation, we have

$$\kappa(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(y_{\alpha}|y_{\alpha}) e^{ikx} dx.$$

Switching the variables I and us, we have

$$\phi(\omega) = \frac{1}{2\sigma} \int_{-\infty}^{\infty} X(gv) e^{i\omega v} dv.$$

Why many when wellow their expenditure as

$$2\pi w(-\omega) = \int_{-\infty}^{\infty} K(\omega)e^{-i\omega t}dt$$

Substituting this equation in eq. (St.28-1), we obtain

$$O(1/\omega) = 2\pi \alpha (-\omega).$$

(b) If in part (a) we have $\sigma(t) = \delta(t+B)$, then we would have $g(t) = K(pt) = e^{-pt}$ and $C(t,\omega) = 2\pi x(-\omega) = 2\pi \delta(-\omega+B) = 2\pi \delta(\omega-B)$.

4.40. When n = 1, $x_1(t) = e^{-it}u(t)$ and $X_1(ju) = 1/(n + ju)$. When n = 2, $x_2(t) = te^{-it}u(t)$ and $X_2(ju) = 1/(n + ju)^2$.

Now, let us insume that the given statement is true when n = m, that in,

$$a_{\infty}(t) = \frac{t^{m-1}}{(m-1)!} e^{-st} u(t) \stackrel{FT}{=} X_{\infty}(y_{n'}) = \frac{1}{(1+y_{n'})^m}$$

a.37. (a) Note that

without

$$\alpha_1(0) = \begin{cases} 1, & \text{ind } \leq \frac{1}{2} \\ 0, & \text{otherwise} \end{cases}$$

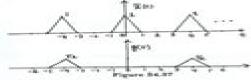
Also, the Fourier transform X_(\$\dot\) of \$1.00 in

$$X_{\mathcal{F}}(p_{m}) = 2^{\min(m/\mathcal{F})}$$

Using the assessitation property we have

$$X(p_{w}) = X_{v}(p_{w})X_{v}(p_{w}) = \left[2^{\min(w)}X_{v}^{T}\right]^{2}$$

(b) The eigend \$00 in as shown in Figure \$4.27



- (x) One possible choice of p(t) is as shown in Figure S4.37
- (66) Note Shed

$$\hat{X}(p\omega) = X(p\omega)\frac{\pi}{2}\sum_{n=1}^{\infty}d(p(\omega - k\frac{\pi}{2})) = G(p\omega)\frac{\pi}{2}\sum_{n=1}^{\infty}d(p(\omega - k\frac{\pi}{2}))$$

This may also be written as

$$\mathcal{R}(p\omega) = \frac{\pi}{2}\sum_{i=1}^{\infty}\mathcal{R}(p\pi k/2)\delta(p\omega - k\frac{\pi}{2}1) = \frac{\pi}{2}\sum_{i=1}^{\infty}\mathcal{Q}(p\pi k/2)\delta(p\omega - k\frac{\pi}{2}1)$$

Clearly, this is possible only if

$$G(j+k/2) = X(j+k/2).$$

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For m = m = 1 we may use the differentiation in frequency property to write,

$$a_{m+1}(t) = \frac{1}{m} a_m(t) \ \delta^{T_m} X_{m+1}(t\omega) = \frac{1}{m} j \frac{dX_m(j\omega)}{d\omega} = \frac{1}{(1+j\omega)^{m+1}}$$

This shows that if we assume that the given statement is true for $n = \infty$, then it is true for $n = \infty + 1$. Since we show this there is a given statement in true n = 0, we in the true tensor in the class it is exceed for n = 2 + 2 = 2, n = 2 + 2 = 4, and so on. Therefore, the given statement

ALASS. DAG Wile Science.

$$\begin{split} g(t) &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{2\pi} \left[X(j\omega) \cdot Y(j\omega) \right] e^{j\omega t} d\omega \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} \frac{1}{2\pi} \left[\int_{-\infty}^{\infty} X(j\theta) Y(j)(\omega - \theta)(d\theta) \right] e^{j\omega t} d\omega \\ &= \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\theta) \left[\frac{1}{2\pi} \int_{-\infty}^{\infty} Y(j)(\omega - \theta)(\omega^{j\omega t}) d\omega \right] d\theta \end{split}$$

(b) Using the frequency shift property of the Fourier transform we have

$$\frac{1}{2m}\int_{-\infty}^{\infty} Y(y) dy - dt(y) e^{2mt} d\omega = e^{2mt} y(y).$$

(a) Conducting the results of parts (a) and (b).

$$g(t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\theta) e^{j\theta t} y(t) d\theta$$

 $= g(t) \frac{1}{2\pi} \int_{-\infty}^{\infty} X(j\theta) e^{j\theta t} d\theta$
 $= g(t) \phi(t).$

4.42. If (i) is a precious signal with Fourier series coefficients a. The fundamental frequency of a(t) is act = 100 rad/sec. From Section 4.2 we know that the Fourier transferre X():v/1 of a(t) is

$$E(p\omega) = \sum_i 2\pi a_i E(\omega - 1000)$$

GAS Since

$$g_{ij}(t) = s(t) \cos(\omega_{ij}t) e^{iT_{ij}} Y_{ij}(p_{ij}) = \frac{1}{2} (X(p(\omega - \omega_{ij})) + X(p(\omega + \omega_{ij})))$$

was basiness

$$Y_{ij}(j\omega) = \pi \sum_{n=1}^{\infty} [a_{nj}\delta(\omega - 100k - \omega_{n}) + a_{nj}\delta(\omega - 100k + \omega_{n})]$$

 $= \pi \sum_{n=1}^{\infty} [a_{nj}\delta(\omega + 100k - \omega_{n}) + a_{nj}\delta(\omega - 100k + \omega_{n})]$ (S4.42-1)

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